

# Ultrasound Sensing with a Photonic Crystal Slab

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**Abstract:** We experimentally demonstrate a compact architecture for ultrasound acoustic sensing based upon a photonic crystal slab (PCS) fabricated on top of a thin mechanically pliable silicon substrate. A noise-equivalent pressure of 2.0 kPa ( $0.072 \text{ kPa}/\sqrt{\text{Hz}}$ ) at 5 MHz is observed. © 2018 The Author(s)

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Acoustic sensing is a technology with myriad applications, ranging from biomedical imaging to non-destructive testing (NDT). The common workhorse for such applications is the piezoelectric transducer [1], which use the piezoelectric effect to both transmit and receive ultrasound signals. However, these devices suffer from narrow bandwidths, and poor sensitivity. Recently, work on capacitive micromachined ultrasound transducers (cMUTs), which are fabricated using CMOS-compatible technology, has helped to improve signal detection bandwidths. However, the detection and generation of ultrasonic signals are still performed electrically. When placed into an array, both piezoelectric and cMUT devices experience significant electrical crosstalk [2], as well as electromagnetic interference (EMI) in general.

One solution to eliminating cross-talk and EMI is to detect the acoustic signals optically. Such implementations often couple the optical and acoustic signals using an optical cavity. The changes in the dimensions of the optical cavity due to the acoustic signal is then imprinted upon the optical reflectivity. Recently, Fabry-Perot devices [3] demonstrated a noise-equivalent pressure of  $90 \mu\text{Pa}/\sqrt{\text{Hz}}$  at 2.5 kHz, but sensitivity at ultrasonic frequencies ( $>1 \text{ MHz}$ ) was not demonstrated. Pi-phase-shifted fiber Bragg gratings (FBG) [4] have been shown to have kPa level sensitivity in the ultrasound regime, but fabricating an FBG sensor array with millimetre spacing remains difficult.

Here, we demonstrate a compact architecture for a photonic crystal slab (PCS) sensor that is capable of detecting ultrasonic acoustic signals. Similar slab sensors were previously used for refractive-index sensing [5]. The device is fabricated using CMOS-compatible technology. A cross-section of the device is shown in Fig. 1a. A 500- $\mu\text{m}$ -thick silicon wafer with thermal oxide and stoichiometric  $\text{Si}_3\text{N}_4$  layer grown on top is used. Circular membranes of diameters 100 to 300 microns is patterned into the silicon layer (back side) using a dry Bosch etch, leaving behind only a 10-micron layer of silicon that deforms under acoustic excitation. A PCS structure is patterned on the nitride layer (front side) of the device through e-beam lithography. The nanohole radius ( $r = 150 \text{ nm}$ ) and pitch ( $a = 1.05 \mu\text{m}$ ) are chosen so that a TE-like Fano resonance peak appears at about 1526 nm (Fig. 1b). The slab structure consists of a  $300 \times 300$  array of nanoholes, centered about the Si membrane. An array of these PCS/membrane combinations are fabricated on a  $15 \text{ mm} \times 15 \text{ mm}$  wafer (Fig. 2a inset), with sensor spacings of 1 mm.

The experimental setup is shown in Fig. 2a. The back side of the structure is immersed in water, while the front side is kept dry using a custom holder. At a distance of 12 cm beneath the silicon membrane is an ultrasound (U/S) transducer (Panametrics C326) with a center frequency of 5 MHz and active element diameter of 9.6 mm. The output of the transducer has been calibrated with an Onda HMB-0200 hydrophone.

Light from a narrow linewidth tunable laser source is used to interrogate the PCS. An  $f = 8 \text{ mm}$  asphere out-

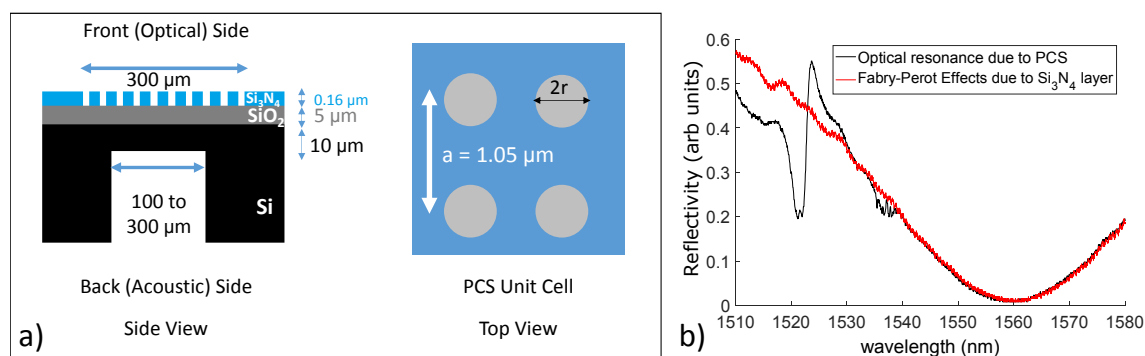


Fig. 1. a) Device geometry. The back side of the sensor consists of a thin membrane of silicon that is mechanically pliable. The front side of the sensor is composed of a photonic crystal slab (PCS) consisting of nanoholes in  $\text{Si}_3\text{N}_4$ . b) The optical reflectivity for the PCS shows a peak at 1526 nm (black trace). Fitting to a Fano lineshape reveals a characteristic width of 2.2 nm, giving us a  $Q$ -factor of 690. The red trace shows the reflectivity from an unpatterned portion of the  $\text{Si}_3\text{N}_4$  layer.

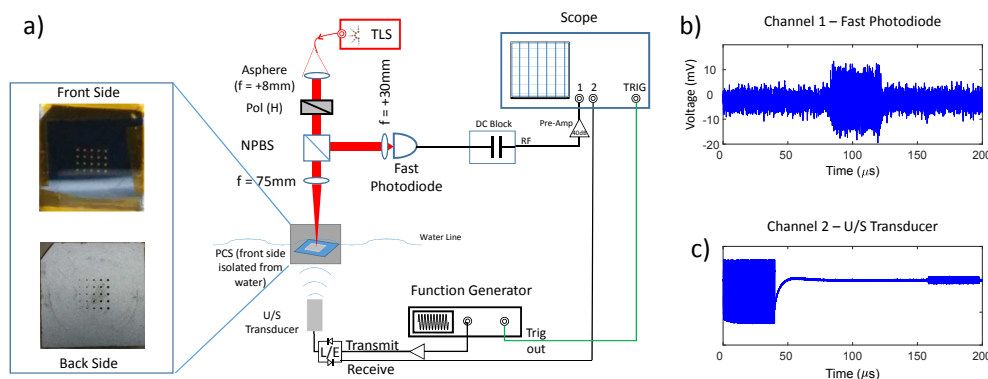


Fig. 2. a) Experimental setup. See text for details. The inset shows the physical chip on which a 2D array of  $4 \times 5$  PCS sensors are present, each separated by 1 mm. On the back side, there are 5 columns of membranes; the leftmost (rightmost) column contains membranes of diameter  $100 \mu\text{m}$  ( $300 \mu\text{m}$ ). b) The optical reflectivity of the PCS is monitored with a fast photodiode and oscilloscope. Notice that the tone burst is observed at  $80 \mu\text{s}$ . c) The transmitted and back-reflected acoustic signals picked up by the ultrasound (U/S) transducer are routed appropriately using an L/E (limiter/expander) circuit. The reflected signal arrives at  $160 \mu\text{s}$ , twice as long as it takes the signal to arrive at the PCS.

couples and collimates the laser light from a fiber, and an  $f = 75 \text{ mm}$  planoconvex lens creates a  $150 \mu\text{m}$  diameter spot size on the slab. The backreflected beam is redirected through a non-polarizing beam splitter (NPBS) to a fast photodiode, whose RF output is then amplified (voltage gain of  $\sim 200$ ), and fed into an oscilloscope.

The ultrasound transducer is driven with 5 MHz sinusoidal tone bursts consisting of 200 periods, at varying voltage amplitudes. The maximum amplitude used allows for the generation of pressure waves with peak-to-peak values of up to 80 kPa at a distance of 12 cm away, and an acoustic beam diameter (FWHM) of 3.7 mm. Figure 2b shows an observed time series as measured by the fast photodiode, after 64 averages (for a total integration time of 2.56 ms). An RMS voltage can be extracted from the trace. Similar traces can be obtained at varying pressures, and the RMS voltage for each can be obtained. For two different membrane diameters (200 and  $300 \mu\text{m}$ ), the applied pressure is varied, and the RMS voltage measured. The results are plotted in Fig. 3a, demonstrating a linear response. The  $300 \mu\text{m}$  membrane is four times as sensitive as the  $200 \mu\text{m}$ , with a noise-equivalent pressure (NEP) of  $2.0 \text{ kPa}$  ( $0.072 \text{ kPa}/\sqrt{\text{Hz}}$ ).

Figure 3b shows the acoustic sensitivity of the  $300 \mu\text{m}$  device as a function of wavelength; the acoustic sensitivity of the PCS correlates well with the slope of the resonance. This measurement acts as a check that the PCS structure contributes primarily to our ability to measure the acoustic waves; the optical response is caused by the mechanical deformation of the membrane and PCS structure in response to the impinging acoustic waves. We anticipate further sensitivity improvements when the front side PCS surface is immersed in water, as pressure-induced index of refraction changes would then be readily observed by the PCS [5], providing another mechanism for acoustic sensing.

In summary, we have demonstrated a compact optical PCS-based ultrasound sensor. It consists of an optical resonator (the PCS) patterned above a mechanically pliable and acoustically sensitive silicon membrane. A noise equivalent pressure of  $0.072 \text{ kPa}/\sqrt{\text{Hz}}$  was observed for the device when detecting 5 MHz ultrasonic waves.

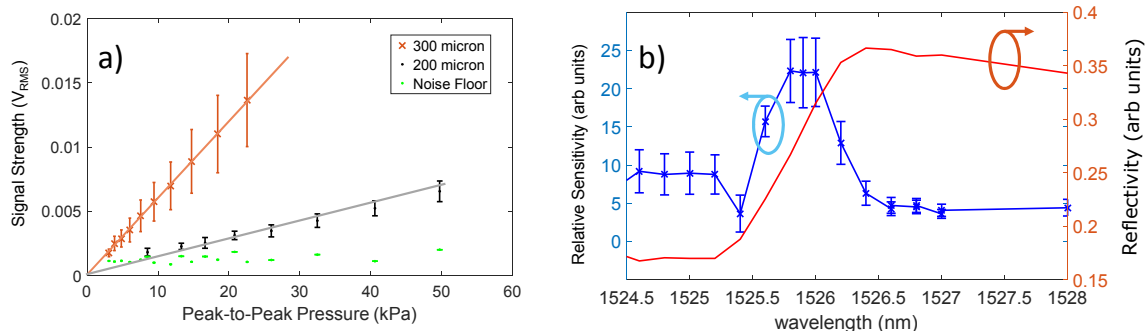


Fig. 3. a) Measured acoustic sensitivity at varying pressures for two different-sized membranes. The  $300 \mu\text{m}$  diameter device yields a four-fold improvement over the  $200 \mu\text{m}$  device. b) The acoustic sensitivity of the sensor is measured as a function of wavelength. The sensitivity (blue curve) is maximized when the slope of the resonance reflectivity (red curve) is maximized.

## References

1. A. Webb and G. C. Kagadis, *Medical Physics* **30**, 2267–2267 (2003).
2. B. Bayram, M. Kupnik, G. G. Yaralioglu, et al., *IEEE transactions on ultrasonics, ferroelectrics, and frequency control* **54** (2007).
3. C. Jan, W. Jo, M. J. Digonnet, and O. Solgaard, et al., *IEEE Photonics Technology Letters* **28**, 123–126 (2016).
4. G. Wissmeyer, D. Soliman, R. Shnaiderman, A. Rosenthal, and V. Ntziachristos, *Opt. Lett.* **41**, 1953–1956 (2016).
5. C. Nicolaou, W. T. Lau, R. Gad, H. Akhavan, R. Schilling, and O. Levi, *Optics express* **21**, 31698–31712 (2013).